

The 2015 outburst of the accreting millisecond pulsar IGR J17511–3057 as seen by *INTEGRAL*, *Swift* and *XMM-Newton*

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ABSTRACT

We report on *INTEGRAL*, *Swift* and *XMM-Newton* observations of IGR J17511–3057 performed during the outburst that occurred between March 23 and April 25, 2015. The source reached a peak flux of $0.7(2) \times 10^{-9}$ erg cm⁻² s⁻¹ and decayed to quiescence in approximately a month. The X-ray spectrum was dominated by a power-law with photon index between 1.6 and 1.8, which we interpreted as thermal Comptonization in an electron cloud with temperature > 20 keV. A broad ($\sigma \simeq 1$ keV) emission line was detected at an energy ($E = 6.9_{-0.3}^{+0.2}$ keV) compatible with the K- α transition of ionized Fe, suggesting an origin in the inner regions of the accretion disk. The outburst flux and spectral properties shown during this outburst were remarkably similar to those observed during the previous accretion event detected from the source in 2009. Coherent pulsations at the pulsar spin period were detected in the *XMM-Newton* and *INTEGRAL* data, at a frequency compatible with the value observed in 2009. Assuming that the source spun up during the 2015 outburst at the same rate observed during the previous outburst, we derive a conservative upper limit on the spin down rate during quiescence of 3.5×10^{-15} Hz s⁻¹. Interpreting this value in terms of electromagnetic spin down yields an upper limit of 3.6×10^{26} G cm³ to the pulsar magnetic dipole (assuming a magnetic inclination angle of 30°). We also report on the detection of five type-I X-ray bursts (three in the *XMM-Newton* data, two in the *INTEGRAL* data), none of which indicated photospheric radius expansion.

Key words. pulsars: general — stars: neutron — X-rays: binaries — X-rays: individual: IGR J17511–3057

1. Introduction

Accreting millisecond pulsars (AMSPs hereafter) are neutron stars (NS) that accrete matter transferred from a low mass ($M_2 \lesssim M_\odot$) companion star (Wijnands & van der Klis 1998); their magnetospheres are able to truncate the disk in-flow and channel the in-falling matter to the regions of the surface close to the magnetic poles, producing coherent pulsations in the X-ray light curve. The extremely quick rotation of AMSPs is attained during a previous Gyr-long phase of sustained mass accretion, and these sources are considered the most immediate progenitors of millisecond radio pulsars (Bisnovatyi-Kogan & Komberg 1974; Alpar et al. 1982; Radhakrishnan & Srinivasan 1982; Archibald et al. 2009; Papitto et al. 2013). So far, accretion driven coherent pulsations have been detected from 17 transient low-mass X-ray binaries (see, e.g. Patruno & Watts 2012, for a review). Pulsations from 15 of these objects have been observed during relatively bright ($L_X \approx \text{few} \times 10^{36}$ erg s⁻¹) X-ray outbursts lasting few-weeks to few-months. In two cases, PSR J1023+0038 and XSS J12270-4859 (Archibald et al.

2015; Papitto et al. 2015), the coherent signal was detected when the source was at a much lower luminosity level ($L_X \approx \text{few} \times 10^{33}$ erg s⁻¹); these two sources were also detected as radio pulsars during X-ray quiescence (Archibald et al. 2009; Roy et al. 2015), similar to the AMSP IGR J18245–2452 (Papitto et al. 2013).

IGR J17511–3057 was discovered by *INTEGRAL* during an outburst in September 2009 (Baldovin et al. 2009; Bozzo et al. 2010). The detection of 4.1 ms coherent pulsations in the *Rossi X-ray Timing Explorer* light curve allowed the identification of the source as an AMSP in a binary system with a 3.47 hr orbital period (Markwardt et al. 2009). The measured pulsar mass function indicated a main-sequence companion star with a mass between 0.15 and 0.44 M_\odot (Papitto et al. 2010). At discovery, 18 type-I X-ray bursts were observed from the pulsar (Altamirano et al. 2010; Bozzo et al. 2010; Papitto et al. 2010; Falanga et al. 2011). Evidence of photospheric radius expansion was not displayed by any of these bursts, and an upper limit to the source distance (6.9 kpc) was provided by Altamirano et al. (2010).

A new outburst of IGR J17511–3057 was detected by *INTEGRAL* on March 23, 2015 (Bozzo et al. 2015a,b).

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Here, we report on a series of *INTEGRAL* and *Swift* observations performed throughout the event, as well as on a *XMM-Newton* Target of Opportunity observation performed three days after the onset of the outburst.

2. Observations

2.1. *Swift*

The *Swift*-XRT (Burrows et al. 2005) observed IGR J17511–3057 from March 23, 2015 at 14:08:39 (UTC) for a total of 19 ks. A log of all available observations for this outburst is reported in Table 1. *Swift*-XRT data, that were collected in both windowed-timing (WT) and photon-counting (PC) mode depending on the source count rate, were processed and analysed using the standard software (FTOOLS v6.17), calibration (CALDB 20150721), and methods. The WT data were never affected by pile-up; on the contrary, the PC data were corrected, when required, by adopting standard procedures (Vaughan et al. 2006). Therefore WT source events were generally extracted from circular regions of 20 pixels (1 pixel $\simeq 2.36''$), while PC source events from annuli with outer radius of 20 pixels and inner radius ranging from 3 to 7 pixels, depending on the severity of pile-up. Background events were extracted from nearby source-free regions. A spectrum was accumulated for each observation and each observing mode, throughout the campaign, with the exception of the last four observations (036–039), for which a single spectrum was extracted to increase statistics. The data were binned to ensure at least 20 counts per energy bin and were fit in the 0.3–10 keV energy range. The latest responses within CALDB were used. Spectral modelling was performed using XSPEC v.12.8.2.

2.2. *XMM-Newton*

XMM-Newton (Jansen et al. 2001) observed IGR J17511–3057 for 76 ks starting on March 26, 2015 at 22:16:52 (UTC). A log of the observations analyzed is given in Table 1. The data were reduced using SAS v.14.0.0.

The EPIC-pn camera was operated in timing mode to achieve a time resolution of 29.5 μ s, and was equipped with a medium optical blocking filter. We removed soft proton flaring episodes characterized by an EPIC-pn 10–12 keV count rate exceeding 0.8 c/s; this reduced the effective exposure to 55.6 ks. In timing mode, the spatial information along one of the optical axis is lost to allow a faster read-out. The maximum number of counts was recorded in pixels characterized by RAWX coordinate 37 and 38. To extract the source emission we considered a 21 pixel wide region (1 pixel $\simeq 4.1''$), spanning from RAWX=27 to 47. Background was extracted far from the source, in a 3 pixel-wide region centered on RAWX=4. Spectra were accumulated using single and double event patterns, and re-binned with the tool SPECGROUP to have at least 25 counts per channel, and not more than three bins per resolution element. For the spectral extraction we followed the recommendations of the latest calibration document on the spectral accuracy of EPIC-pn in fast modes (Smith et al. 2015¹) and applied the special gain, together with the rate dependent CTI corrections.

Three type-I X-ray bursts were detected during the *XMM-Newton* observation. To analyse the *persistent* emission we created good time intervals (GTIs) that eliminated a time interval starting 10 s before and ending 150 s after each of these bursts. The burst onset was identified as the first 1 s-long bin of the light curve that exceeded by more than 5σ the average *persistent* countrate of 65.6 c/s. Apart from the bursts, no evident variability trend was seen over time scales going from a few seconds to the length of the observation. At the average *persistent* count rate observed photon pile up is not expected to affect significantly the spectral response of the EPIC-pn in timing mode (Smith et al. 2015).

The Reflection Grating Spectrometers (RGS) was operated in standard spectroscopy mode. We extracted spectra from the first order of diffraction, in which a count rate of 5.0 and 5.9 c/s was observed by the RGS1 and RGS2 cameras, respectively. We removed the same time intervals characterized by high flaring background and type-I X-ray bursts as it was done for the EPIC-pn.

The EPIC MOS1 and MOS2 cameras operated in Large Window and Timing mode, respectively. We did not consider these data in the analysis presented here as $\simeq 21\%$ and $\simeq 12\%$ of the photons detected by the two cameras respectively suffered from pile-up distortion of the spectral response, as evaluated with the EPATPLOT tool.

2.3. *INTEGRAL*

IGR J17511–3057 was observed within the field of view (FoV) of the instruments on-board *INTEGRAL* from satellite revolution 1517 (starting on 2015 March 3) to 1533 (ending on 2015 April 25).

We analyzed all the publicly available *INTEGRAL* data and the data for which we obtained data rights on the source during *INTEGRAL* AO12 using version 10.1 of the Off-line Scientific Analysis software (OSA) distributed by the *INTEGRAL* Science Data Center (ISDC; Courvoisier et al. 2003). The *INTEGRAL* data on the source were accumulated during observations of the Galactic Center, which are executed following a rectangular pattern (5x5 pointings, or science windows, SCWs) around the Galactic Center position, with typical duratins of 2–3 ks and a 2.17 degree step between successive pointings. Only SCWs in which the source was located within 4.5° from the center of the JEM-X FoV (Lund et al. 2003) were included in the analysis of the data from this instrument. For IBIS/ISGRI, we included all SCWs where the source was located within 12° from aim point of the instrument in order to avoid calibration uncertainties (as suggested in the OSA manual). We first extracted both the IBIS/ISGRI (Ubertini et al. 2003; Lebrun et al. 2003) mosaics in the 20–80 keV and 80–100 keV energy band and noticed that there was no significant detection of the source in the hard band. IGR J17511–3057 was significantly detected in both the JEM-X mosaics extracted in the 3–10 keV and 10–20 keV energy bands. The ISGRI and JEM-X lightcurves of the source were extracted with a time resolution of 1 SCW in the 20–80 keV and 3–20 keV, respectively and later rebinned, to improve the statistics, with a time resolution of 2 hours (see panels (c) and (d) in Fig. 1).

We accumulated a single spectrum of the source for each of the two JEM-X units and for IBIS/ISGRI using the data

¹ <http://xmm2.esac.esa.int/docs/documents/CAL-TN-0083.pdf> in revolution 1522, between 2015 March 25 at 03:55:44, and

Table 1. Log of the *Swift* and *XMM-Newton* observations considered in this paper

Sequence	Obs/Mode	Start time (UT) (yyyy-mm-dd hh:mm:ss)	End time (UT) (yyyy-mm-dd hh:mm:ss)	Exposure (s)
<i>Swift</i>				
00031492020	XRT/WT	2015-03-23 14:08:39	2015-03-23 16:18:57	860
00031492021	XRT/WT	2015-03-24 16:00:13	2015-03-24 16:15:58	926
00031492022	XRT/PC	2015-03-25 01:48:43	2015-03-25 06:39:55	923
00031492024	XRT/WT	2015-03-28 15:54:29	2015-03-28 16:05:58	682
00031492025	XRT/WT	2015-03-28 19:16:09	2015-03-28 19:20:58	266
00031492026	XRT/WT	2015-03-29 20:30:24	2015-03-29 20:46:58	981
00031492027	XRT/WT	2015-04-02 04:30:10	2015-04-02 06:11:26	97
00031492027	XRT/PC	2015-04-02 04:31:20	2015-04-02 06:20:54	1382
00031492028	XRT/PC	2015-04-04 13:33:42	2015-04-04 13:38:28	266
00031492029	XRT/WT	2015-04-06 13:52:14	2015-04-06 15:11:08	186
00031492029	XRT/PC	2015-04-06 13:52:49	2015-04-06 15:16:55	1241
00031492031	XRT/PC	2015-04-07 13:41:39	2015-04-07 13:57:55	973
00031492032	XRT/PC	2015-04-08 02:35:18	2015-04-08 02:50:56	920
00031492033	XRT/PC	2015-04-08 10:11:57	2015-04-08 11:58:54	1449
00031492034	XRT/PC	2015-04-12 00:58:10	2015-04-12 05:45:56	1316
00031492035	XRT/PC	2015-04-14 09:55:08	2015-04-14 10:19:55	1487
00031492036	XRT/PC	2015-04-16 10:00:40	2015-04-16 10:24:55	1449
00031492037	XRT/PC	2015-04-18 21:23:41	2015-04-18 23:04:55	757
00031492038	XRT/PC	2015-04-22 09:53:57	2015-04-22 10:06:54	777
00031492039	XRT/PC	2015-04-25 04:47:13	2015-04-25 17:56:54	2038
<i>XMM-Newton</i>				
0770580301	EPIC-pn	2015-03-26 22:51:47	2015-03-27 19:11:33	73186
	RGS1	2015-03-26 22:16:52	2015-03-27 19:07:38	75045
	RGS2	2015-03-26 22:17:00	2015-03-27 19:08:29	75089

March 26 at 23:14:27, leading to an effective on-source exposure time of 31.2 ks and 24.3 ks, respectively. Such a time interval was chosen because it is the closest to the *XMM-Newton* observation, having an overlap of 4.5 ks with it. In particular, we selected for the extraction of this spectrum only the SCWs where the source was detected at more than 5σ and within an off-axis angle of 4.5° in order to optimize the signal-to-noise ratio and have a fully simultaneous JEM-X+ISGRI spectrum. The spectra response matrices were generated using the standard energy binning of 16 channels for JEM-X and 13 channels for IBIS/ISGRI. The fit to these spectra was carried out simultaneously with the *XMM-Newton* spectrum and discussed in Sec. 2.2.

We also extracted the JEM-X1 and JEM-X2 lightcurves with a time resolution of 2 s to search for type-I X-ray bursts, and identified two, none in the time window overlapping with *XMM-Newton* exposure. The properties of these bursts are discussed in Sec. 3.3. On the other hand, we could not identify any burst in the ISGRI light curve.

3. Results

3.1. The persistent emission

The *Swift*-XRT 0.3–10 keV spectra were modelled simultaneously with an absorbed power-law, forcing the absorption column to take the same value in every observation. This resulted in an estimate of the absorption column density (described with the model PHABS in XSPEC)

of $N_H = 1.18(4) \times 10^{22} \text{ cm}^{-2}$. The chi squared of the fit was 1632 for 1411 degrees of freedom. The probability of obtaining a fit chi-squared as high or higher if the data are drawn from the model distribution is 3.5×10^{-5} , smaller than the significance level of 2.7×10^{-3} we set to determine the goodness of the fit. A probability always larger than the significance level adopted. Letting the absorption column as a free parameter in different observations did not yield a statistically significant decrease of the fit chi-squared. However, we did not find any evident unmodelled residuals that motivated the addition of more spectral components, and the chi squared of the individual spectra resulted in a null hypothesis probability always larger than the significance level adopted. Systematic errors we did not take into account possibly yielded the relatively large residuals of the simultaneous modelling of the *Swift*-XRT spectra with an absorbed power law. The 0.5–10 keV unabsorbed flux attained a maximum of $0.7(2) \times 10^{-9} \text{ erg cm}^{-2} \text{ s}^{-1}$ on March 24, during observation 00031492021 (see Table 1). The flux decay started around March 29 and was characterized by a roughly linear decay, until the flux decreased abruptly on April 16 (i.e. 25 days after the first detection of the outburst) to a level $\simeq 100$ fainter than the outburst peak ($F(0.5 - 10 \text{ keV}) = 7.2(8) \times 10^{-12} \text{ erg cm}^{-2} \text{ s}^{-1}$ measured coadding the observations performed between April 16 and April 25). The photon index of the power-law that fitted the *Swift*-XRT spectra laid between 1.3 and 1.6, becoming much steeper ($\Gamma = 2.1(2)$) in the last group of observa-

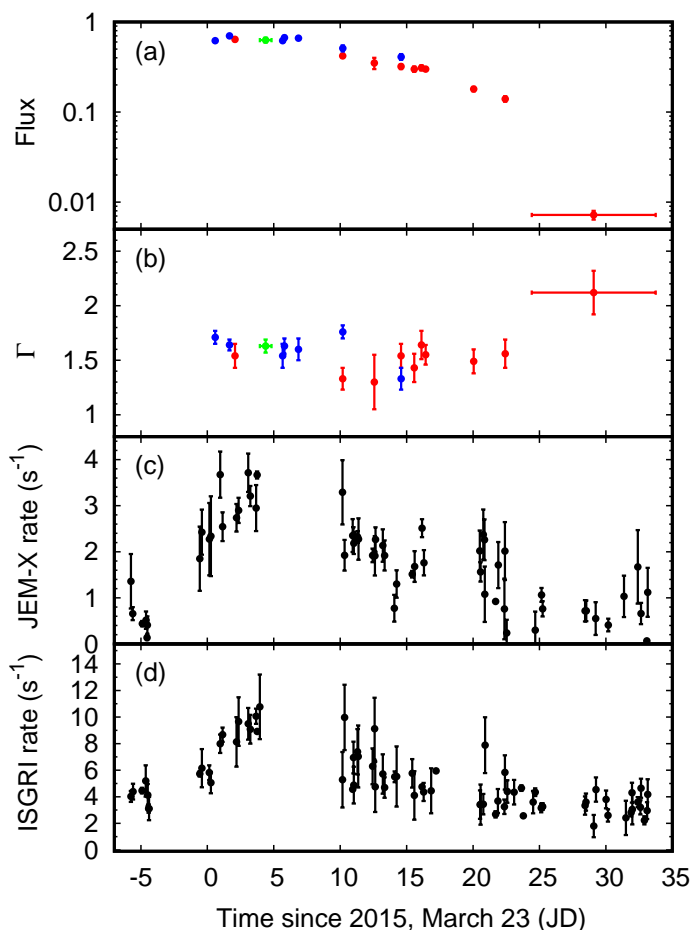


Fig. 1. 0.5–10 keV unabsorbed flux (in units of 10^{-9} erg cm $^{-2}$ s $^{-1}$; panel a) and photon index of the best fitting power-law (panel b) measured by the *Swift*-XRT in WT (blue points) and PC mode (red points), and by *XMM-Newton* (green point). JEM-X (panel c) and ISGRI (panel d) light curves were binned with a time resolution of 2 hours.

tions performed when the source had already faded significantly. The values of the unabsorbed 0.5–10 keV flux and power law index are plotted in the upper panels (a and b) of Fig. 1 using blue and red points for observations performed in WT and PC modes, respectively.

The spectra observed on March 26, 2015 by the EPIC-pn (1.4–11 keV)² and the two RGS cameras (0.6–2.0 keV) on-board *XMM-Newton* were modelled simultaneously by treating the normalization of the two RGS spectra with respect to the EPIC-pn as free parameters (see top panel of Fig. 2). We used a thermal Comptonization model (NTHCOMP, Zdziarski et al. 1996; Życki et al. 1999) modified by the interstellar absorption. No high-energy cut-off was detected in the *XMM-Newton* energy band and we fixed the temperature of the Comptonizing electron cloud to $kT_e = 30$ keV, in line with the results obtained modelling the INTEGRAL spectrum observed immediately before the XMM pointing (see below). The Comptonization component was described by an asymptotic photon index $\Gamma = 1.62^{+0.03}_{-0.11}$, and a tem-

perature of soft input photons of $kT_s < 0.98$ keV, yielding a high value of the fit chi-squared of 1963.3 for 1286 degrees of freedom. To account for the residuals at soft ($kT \lesssim 1.5$ keV) energies we added a single-temperature black-body component (BBODYRAD in XSPEC) with $kT_{bb} = 0.59^{+0.15}_{-0.11}$ keV, and a disk emission component (DISKBB in XSPEC) with temperature of $kT_{in} = 0.26^{+0.04}_{-0.03}$ keV. The addition of these two components yielded a significant decrease of the χ^2 by $\Delta\chi^2 = -74.3$ (giving $\chi^2 = 1888.9$ for 1284 d.o.f.) and -124.6 (giving $\chi^2 = 1764.29$ for 1282 d.o.f.), for the addition of two degrees of freedom each, respectively. We evaluated with an F-test that such improvements of the model χ^2 are equivalent to a probability of $< 10^{-11}$ that they are due to chance. The normalizations of these two thermal components indicate apparent radii of $R_{in}/(\cos i)^{1/2} = (34 \pm 20) d_{6.9}$ km and $R_{bb} = 3.2^{+4.6}_{-1.0} d_{6.9}$ km for the disk and the blackbody component, respectively. Here, $d_{6.9}$ is the distance to the source in units of 6.9 kpc. Local residuals of the EPIC-pn data at energies of 1.8, 2.2 and 6.9 keV were modelled using three Gaussian emission features. The first two features are narrow, and their energies are compatible with calibration residuals of the Si K and Au M edges, known to frequently affect the EPIC-pn spectra in timing mode. The feature at higher energy had a centroid at $E = 6.9^{+0.2}_{-0.3}$ keV, compatible with the K α transition of ionized Fe. With the inclusion of this feature the model χ^2 decreased by $\Delta\chi^2 = 130.8$ for the addition of three free parameters (see Protassov et al. 2002 and Stewart 2009 for a discussion of the use of the fit chi-squared to measure the significance of spectral lines). The residuals obtained with and without the addition of the Fe emission line to the model are shown in the bottom and middle panel of Fig. 2, respectively. The line width ($\sigma = 1.1^{+0.4}_{-0.2}$ keV) and strength (equivalent width of $\simeq 150$ eV), suggested broadening in the inner parts of an accretion disk. An absorption edge at an energy compatible with the OVIII transition (0.871 keV) was also added to model residuals of the RGS spectra. With the addition of such emission features, the reduced χ^2 of the fit attained a value of 1441.8 (for 1274 degrees of freedom). The probability that a value of the χ^2 as high or higher is produced if the data are drawn from the model distribution is 6.8×10^{-4} . This probability is smaller than significance level we set ($p = 2.7 \times 10^{-3}$), but the absence of evident residuals, and the 2% uncertainty quoted by Smith et al. (2015)³ for the relative effective area calibration of the EPIC-pn, motivated us to accept this model without the addition of further components. Model parameters of the best-fit are listed in the central column of Tab. 2.

The broadness of the Fe line motivated us to attempt modelling it using a DISKLINE component (Fabian et al. 1989). However, the best fit was obtained only for extreme values of the parameters that control the line width ($R_{in} < 8.5 R_g$, $i > 67^\circ$), and did not yield a significant improvement of the description of the spectrum, as the chi-squared of the fit decreased by $\Delta\chi^2 = 3.7$ for the addition of three free parameters.

We then fitted the *XMM-Newton* EPIC-pn and RGS spectra together with the *INTEGRAL* JEM-X (5–25 keV) and ISGRI (20–150 keV) spectra which were taken on March, 26 (between 03:55:44 and 23:14:27). An AMSP like IGR J17511–3057 does not usually show drastic spectral

² EPIC-pn data below 1.4 keV were discarded to avoid the spurious soft excess that sometimes appears in the EPIC-pn spectra obtained in fast modes, and of which we found hints also in this observation (see, e.g. Hiemstra et al. 2011).

³ <http://xmm2.esac.esa.int/docs/documents/CAL-TN-0018.pdf>

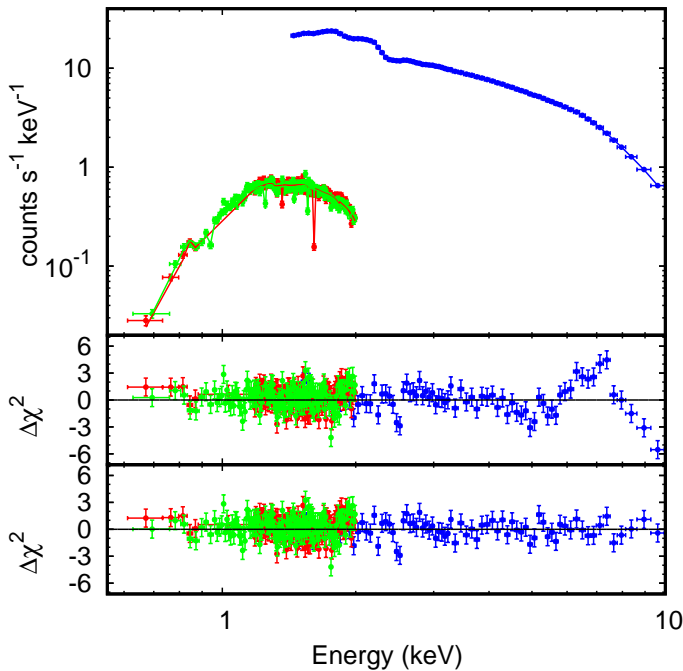


Fig. 2. Top panel shows the X-ray spectrum of IGR J17511–3057 observed by the EPIC-pn (blue points), RGS1 (red points), and RGS2 (green points) during the *XMM-Newton* observation. Spectra were rebinned for graphical purposes. The bottom and middle panels show residuals with respect to the model listed in the central column of Table 2 (see text for details), with and without the addition of an emission line centered at $E_1 = 6.9^{+0.2}_{-0.3}$ keV to the model, respectively.

variability on timescales of a day or less (Ibragimov et al. 2011), and the analysis of Swift data presented earlier where no significant spectral changes were detected over the outburst, supports the simultaneous modelling of the *XMM-Newton* and *INTEGRAL* data in spite of the limited overlap ($\simeq 4.5$ ks) between these observations. The ISGRI spectrum allowed us to weakly constrain the electron temperature of the Comptonized component as $kT_e = 33^{+69}_{-11}$ keV, while the rest of the spectral parameters are basically unchanged with respect to those determined by *XMM-Newton* spectra alone. The spectrum and residuals from the best-fit are plotted in Fig. 3, while model parameters are listed in the rightmost column of Tab. 2.

We also considered the addition to the model of a component describing the reflection of the Comptonizing photon spectrum off the accretion disk. We removed the Gaussian describing the Fe emission line, and added a REFLIONX (Ross & Fabian 2005) component to the model, convolved with relativistic disk blurring (RDBLUR, Fabian et al. 1989). The parameters of the disk reflecting surface were held fixed to inner and outer radii of $R_{in} = 10 R_g$ and $R_{out} = 1000 R_g$, respectively, and index of the power law dependence of emissivity $\beta = -2$. The best fit was obtained for a ionization parameter $\xi = 1500^{+1000}_{-500}$, and a flux of $0.10(5) \times 10^{-10}$ erg cm $^{-2}$ s $^{-1}$ in the reflection component, equal to $\simeq 10\%$ of the irradiating flux and similar to the value observed by Papitto et al. (2011). The addition of reflection did not yield a significant decrease of the model χ^2 ($\chi^2 = 1457.1$ for 1291 degrees of freedom), though, and we could not conclude that a reflection component was signifi-

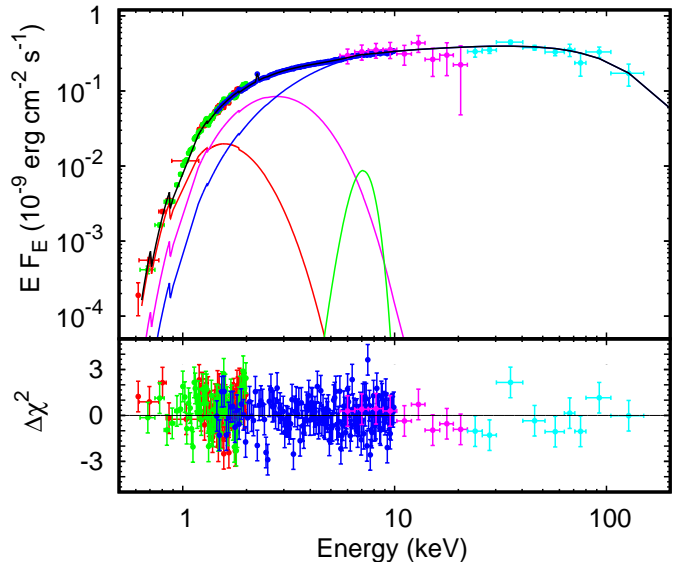


Fig. 3. Top panel shows the X-ray spectrum of IGR J17511–3057 observed by the EPIC-pn (blue points), RGS1 (red points), RGS2 (green points), JEM-X (magenta points) and ISGRI (cyan points) during the 2015 outburst. RGS1 and RGS2 spectra were rebinned for graphical purposes. The best-fit model (black dashed line), the Fe line (green dashed line), the disk (red dashed line), single-temperature black body (magenta dashed line), and the Comptonized (blue dashed line) are also shown. Bottom panel shows residuals with respect to the best fitting model.

cantly detected in the considered dataset, probably because of low counting statistics.

3.2. Timing analysis

To perform a timing analysis of the coherent signal shown by the source in the *XMM-Newton* EPIC-pn data, we first converted the times of arrival of the X-ray photons to the Solar system barycenter using the source position determined by Nowak et al. 2009 using Chandra observations (see also Paizis et al. 2012). Coherent pulsations at 244.8 Hz were easily detected at an rms amplitude of $\simeq 15$ per cent in 50 s-long intervals. Starting from the orbital solution provided by Riggio et al. (2011b), we applied standard timing techniques (see, e.g. Papitto et al. 2011) to study the evolution of the pulse phase computed over 500 s-long intervals. The values of the orbital parameters we obtained, namely the semi-major axis of the NS orbit $a \sin i/c$, the orbital period P_{orb} , the epoch of passage at the ascending node T^* and eccentricity e , as well as the value of the spin frequency ν , and spin frequency derivative $\dot{\nu}$, are given in Table 3. The spin frequency derivative could not be constrained over the relatively short exposure of the *XMM-Newton* observation. As we show in Sec. 4, spin up torques are expected to increase the pulsar frequency at a rate of $\approx \text{few} \times 10^{-13}$ Hz s $^{-1}$, two orders of magnitude lower with respect to the upper limit we measured ($|\dot{\nu}| < 1.2 \times 10^{-11}$ Hz s $^{-1}$, at 3σ confidence level). Because of the relatively short exposure, the best fit values of the spin frequency correlated to a high degree with the (unconstrained) spin frequency derivative. Similarly to what was done by Papitto et al. (2010) for the 2009 outburst of IGR J17511–3057, we quote in Tab. 3 the pulsar spin frequency obtained fixing $\dot{\nu} = 0$ in the fit (i.e. the average spin fre-

Table 2. Spectral parameters of IGR J17511–3057

Parameter	<i>XMM</i>	<i>XMM-IGR</i>
N_H (10^{22} cm $^{-2}$)	$1.03^{+0.10}_{-0.07}$	0.99(5)
τ_{OVIII}	0.18(7)	
kT_{in} (keV)	$0.26^{+0.04}_{-0.03}$	0.17(7)
$R_{in}/(\cos i)^{1/2}$ ($d_{6.9}$ km)	34 ± 20	20^{+5}_{-7}
kT_{bb} (keV)	$0.59^{+0.16}_{-0.12}$	0.62(3)
R_{bb} ($d_{6.9}$ km)	$3.2^{+4.6}_{-1.0}$	$6.6^{+0.7}_{-0.9}$
Γ	$1.63^{+0.03}_{-0.11}$	$1.85^{+0.22}_{-0.15}$
kT_s (keV)	< 0.98	1.3(2)
kT_e (keV)	(30.0)	33^{+69}_{-11}
F_{-9}^{nth} (erg cm $^{-2}$ s $^{-1}$)	0.49 ± 0.15	1.18(6)
E_1 (keV)	$6.89^{+0.18}_{-0.27}$	$6.93^{+0.23}_{-0.21}$
σ_1 (keV)	$1.07^{+0.40}_{-0.22}$	$0.78^{+0.4}_{-0.3}$
N_1 (10^{-4} cm $^{-2}$ s $^{-1}$)	$5.7^{+9.6}_{-1.9}$	$2.2^{+2.2}_{-1.1}$
E_2 (keV)	1.86(2)	1.86(2)
N_2 (10^{-4} cm $^{-2}$ s $^{-1}$)	1.2(4)	1.2(4)
E_3 (keV)	2.24(1)	2.24(7)
N_3 (10^{-4} cm $^{-2}$ s $^{-1}$)	2.7(6)	2.7(3)
RGS1/EPN	0.98(1)	0.98(1)
RGS2/EPN	0.98(1)	0.98(1)
JEM-X/EPN	...	0.58(9)
ISGRI/EPN	...	$1.13^{+0.39}_{-0.25}$
F_{-9} (erg cm $^{-2}$ s $^{-1}$)	0.63(4)	1.43(6)
$\chi^2(\text{dof})$	1441(1274)	1457(1291)

Notes. Best fit parameters of the spectrum observed by *XMM-Newton* (RGS in the 0.6–2.0 keV energy band and EPIC-pn in the 1.4–10 keV range) on March 26, 2016 (central column), and of the simultaneous spectrum observed by *XMM-Newton* and INTEGRAL (JEM-X in the 5–25 keV, and ISGRI in the 20–150 keV band; rightmost column). Fluxes are unabsorbed, given in units of 10^{-9} erg cm $^{-2}$ s $^{-1}$, and evaluated in the 0.5–10 keV for the *XMM* spectrum, and the 0.5–100 keV energy band for the *XMM-IGR* spectrum, respectively. Uncertainties are quoted at the 90% confidence level.

quency over the considered interval). A compatible value of the spin frequency is obtained if a value of the spin-up rate of the order of that expected from the observed X-ray luminosity ($\dot{\nu} = 2 \times 10^{-13}$ Hz s $^{-1}$) is considered, justifying our choice.

The error on the position of the source is an additional source of uncertainty on the value of the spin frequency (see, e.g. Lyne & Graham-Smith 1990). To evaluate the frequency shift produced by a difference between the actual value of the ecliptic coordinates of the source and that used to report X-ray photons to the Solar system barycenter, we used (see, e.g., Eq. 4 in Papitto et al. 2011):

$$\delta\nu_{pos} = \nu y \frac{2\pi}{P_{\oplus}} [\cos M_i \cos \beta \delta\lambda + \sin M_i \sin \beta \delta\beta]. \quad (1)$$

Here, y is the Earth distance from the Solar System barycentre in lt-s, λ and β are the ecliptic longitude and latitude, respectively, $\delta\lambda$ and $\delta\beta$ are the respective uncertainties, $M_i = [2\pi(T_i - T_{\gamma})/P_{\oplus}] - \lambda$, T_i is the start time of observations considered, T_{γ} is the nearest epoch of passage

Table 3. Spin and orbital parameters of IGR J17511–3057

$\langle \nu \rangle$ (Hz)	244.83395112(3)
$\dot{\nu}$ (Hz s $^{-1}$)	$< 1.2 \times 10^{-11}$
$a \sin i/c$ (lt-s)	0.275196(4)
P_{orb} (s)	12487.53(2)
T^* (MJD)	57107.8588090(15)
e	$< 8 \times 10^{-5}$
$\chi^2(\text{d.o.f.})$	141.5(121)

Notes. The timing solution is referred to the epoch $T_0 = 57107.954157$ MJD. The error on the spin frequency does not take into account the uncertainty driven by the indetermination on the position of the source ($\sigma_{\nu}^{pos} = 3 \times 10^{-8}$ Hz)

at the vernal point, and P_{\oplus} is the Earth orbital period. The error of the position of the source ($0.6''$ at 90% confidence level, Nowak et al. 2009; Paizis et al. 2012, which translates into $\sigma_{\lambda} = 1.5 \times 10^{-6}$ rad, $\sigma_{\beta} = 1.8 \times 10^{-6}$ rad) yielded an uncertainty of $\sigma_{\nu}^{pos} = 3 \times 10^{-8}$ Hz on the value of the measured spin frequency. Adding in quadrature this uncertainty to the statistical error gave an error of $\sigma_{\nu} = 4 \times 10^{-8}$ Hz on the average spin frequency measured by *XMM-Newton* in 2015.

Folding the entire EPIC-pn observation around the best fitting value of the spin frequency, we obtained the 0.3–10 keV pulse profile displayed in Fig. 4. It was successfully modelled by four harmonic components, similar to the pulse profile observed in 2009 (Papitto et al. 2010).

We searched for coherent pulsations in the hard X-rays ISGRI band, considering the data observed by ISGRI from 2015 March 22 at 08:56:29 to April 9 at 05:52:01 UTC. We converted the arrival times of the events recorded by ISGRI at the Solar system barycenter using the position derived by Chandra, and then to the line of nodes using the ephemeris listed in of Table 3. We employed a customary software dedicated to phase-resolved spectroscopy of X-ray binaries with IBIS (Segreto & Ferrigno 2007) to extract the background subtracted pulse profile in the broad 20–100 keV energy band. We detected the signal at the pulsar spin frequency determined by *XMM-Newton*, but the low signal to noise ratio of the ISGRI data prevented us from obtaining an independent measurement. The pulse profile is displayed in Fig. 5. Fitting the profile with a constant and a sinusoid, we determined that the ratio between the sine amplitude and the constant as $14 \pm 3\%$, in agreement with what was found by Falanga et al. (2011) during the 2009 outburst of this source. Splitting the energy range in two bands, we obtained amplitudes of $22 \pm 6\%$ and $12 \pm 3\%$ for the energy ranges 20–30 keV and 30–100 keV, respectively, suggesting a decrease of the pulsed emission at higher energy.

In order to study the aperiodic time variability, we produced a power density spectrum of the 0.5–10 keV EPIC-pn time series with 59 μ s time resolution (yielding a Nyquist frequency of 8468 Hz). We averaged the Leahy-normalized fast Fourier transforms performed over approximately 8 s-long intervals (for a total of 9457 spectra, each extending down to 0.1 Hz). The resulting average power density spectrum was rebinned as a geometrical series with a ratio of 1.02. The 0.1–500 Hz spectrum was modelled with the sum of two flat-top noise components modeled with Lorentzian functions centered at $\nu = 0$ with widths

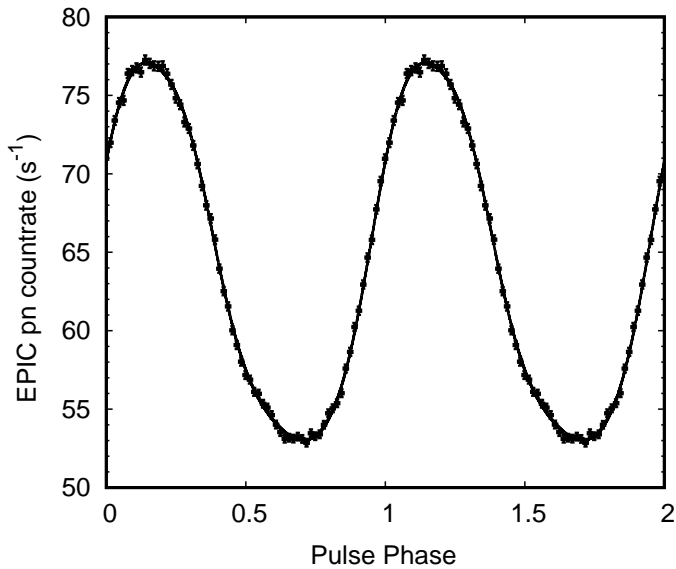


Fig. 4. 0.3–10 keV pulse profile of IGR J17511–3057 as observed by the EPIC pn during the 2015 outburst. The best fitting model (solid line) is composed by four harmonics with background-subtracted, rms amplitudes of 14.31(6)%, 1.53(6)%, 0.72(6)% and 0.26(6)%, respectively. Two cycles are displayed for clarity.

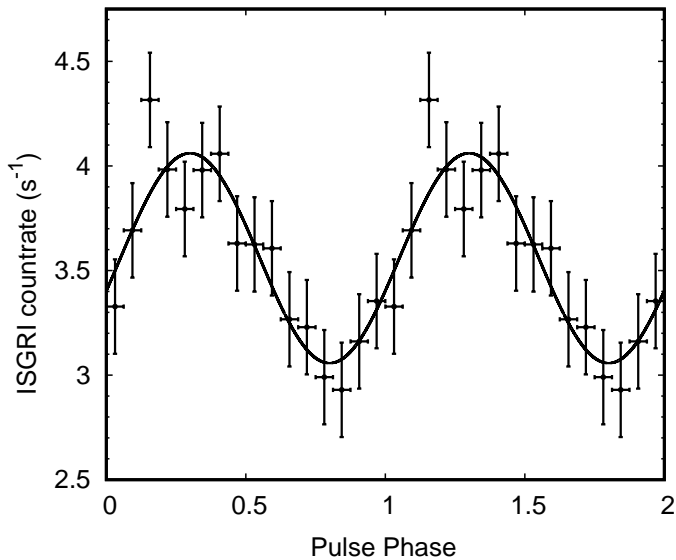


Fig. 5. Background subtracted pulse profile of IGR J17511–3057 in the 20–100 keV energy band extracted from 221 ks of dead-time corrected on-axis equivalent exposure time in the IBIS/ISGRI data (from 2015-03-22 08:56:29 to 04-09 05:52:01 UTC); the vertical axis displays equivalent count per second for on-axis pointing.

$W_1 = (3 \pm 1) \times 10^{-2}$ Hz and $W_2 = 11.5 + / - 1.3$ Hz, and a discrete feature centered at $\nu_3 = 0.5 \pm 0.1$ Hz with width $W_3 = 1.2 \pm 0.2$ Hz, giving a fit χ^2 of 611.2 for 522 degrees of freedom. We searched for kHz QPOs fitting a Lorentzian with quality factor fixed to a value of 4, but found no significant QPO with a 3σ upper limit of 1% on the rms amplitude, a value lower than that characterizing the high frequency QPOs reported by Kalamkar et al. (2011). The average power spectrum and residuals with respect to the best fit model are plotted in Fig. 6.

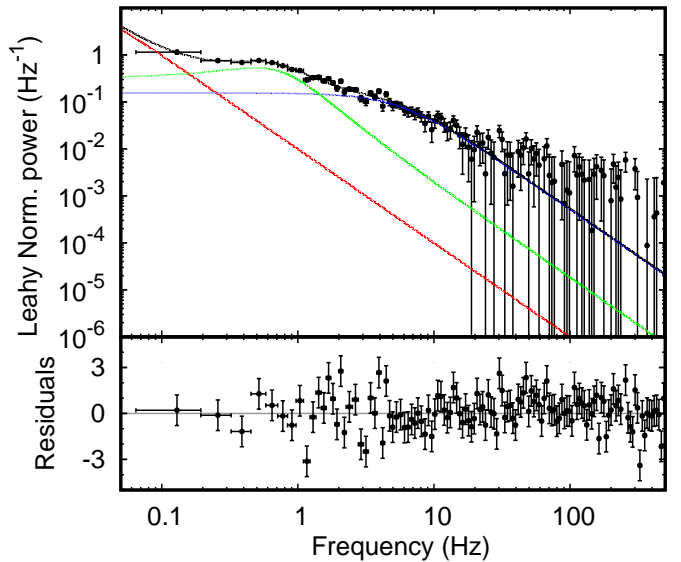


Fig. 6. Average Leahy normalized power spectrum of the EPIC-pn time series (top panel). A counting noise constant level of 1.9993(4) is subtracted to the power. The two flat-top noise components are plotted with a red and a blue solid line, respectively, and the discrete feature centered at $\nu_3 = 0.5 \pm 0.1$ Hz as a green solid line (see text for details). The spike at the pulsar spin frequency was removed for graphical purposes. Bottom panel shows residuals in units of the standard deviation σ .

3.3. Type-I X-ray bursts

During the *XMM-Newton* observation, three type-I X-ray bursts were detected, sharing very similar properties. The light curve of the first burst is displayed in panel (a) of Fig. 7. The start times of the three bursts were $T_1 = 57108.01023(1)$ MJD, $T_2 = 57108.34130(1)$ MJD, and $T_3 = 57108.67847(1)$ MJD, respectively. The burst rise times were ≈ 1 s, and were followed by exponential decays which started a few seconds after the maxima. The measured decay e-folding times were $\tau_1 = 12.1(2)$ s, $\tau_2 = 12.6(2)$ s and $\tau_3 = 12.4(2)$ s for the three bursts, respectively. The time elapsed between consecutive bursts was $\Delta t_1 = 28.6$ ks and $\Delta t_2 = 29.1$ ks.

We extracted the EPIC-pn burst spectra removing the two brightest pixel columns to minimize the effect of pile up, and subtracted the spectrum of the persistent emission as the background. Close to the burst peak spectra were extracted over 2 s long intervals; at later times the exposure was increased to ensure that the counting statistics allowed a meaningful spectral modelling. Spectral channels were rebinned in the same way as the persistent spectrum. The spectra extracted over the different intervals were simultaneously fit by an absorbed blackbody with variable temperature and radius, resulting in a fit χ^2 of 510.9 (for 556 degrees of freedom), 559.43 (522) and 589.16 (518) for the three bursts, respectively. The 0.5–10 keV burst flux attained a maximum value of $(1.5 \pm 0.1) \times 10^{-8}$ erg cm $^{-2}$ s $^{-1}$. The temperature decreased steadily from ≈ 3 keV during the burst evolution (see panel b of Fig. 7), while the apparent radius of the emitting region peaked at $\approx 6 d_{6.9}$ km after ten seconds since the burst onset, and then decayed smoothly (see panel c of Fig. 7).

Coherent oscillations at the pulsar spin frequency were detected at a somewhat lower rms amplitude (≈ 10 per

cent) than that of persistent pulsations ; see panel d of Fig. 7). Considering that the burst emission was up to 25 times brighter than the persistent flux, such a decrease of the pulse amplitude could not be ascribed to the increase of the unpulsed flux, but suggests instead that the oscillations detected during the burst were different than the coherent signal observed during quiescence. Pulsations were detected both close to the burst peak and during the burst decay. The phase of burst oscillations measured over 5-s long intervals is generally consistent within the uncertainties with the phase of the persistent profile, even if deviations up to 0.2 in phase are observed close to the burst peaks.

Two bursts were detected by *INTEGRAL*, with trigger times of $T_4=57116.43276(2)$ MJD and 57120.11919(2) MJD. From the JEM-X lightcurves we measured a peak count-rates of 344.6 ± 47.6 c/s (3–20 keV) and 177.5 ± 38.8 c/s (3–20 keV), respectively for the two bursts. Because JEM-X spectra cannot be extracted for time intervals shorter than 8 s, we estimated the peak flux by comparing the above count-rates with that of the Crab in the same energy band as obtained from the *INTEGRAL* calibration observations performed in the satellite revolution 1520 (on 2015 March 19). We obtained a 3–20 keV X-ray flux of 1.3 ± 0.2 Crab and 0.7 ± 0.1 Crab, respectively. These translate into $\sim 3.0 \times 10^{-8}$ erg cm $^{-2}$ s $^{-1}$ and $\sim 1.6 \times 10^{-8}$ erg cm $^{-2}$ s $^{-1}$. The e-folding decay times of the two bursts measured from the JEM-X lightcurves were 7.3 ± 0.1 s and 6.0 ± 0.1 s.

4. Discussion

During the outburst detected in early 2015, IGR J17511–3057 attained a peak flux of $\simeq 0.6\text{--}0.7 \times 10^{-9}$ erg cm $^{-2}$ s $^{-1}$ (0.5–10 keV, unabsorbed, see Sec. 3.1). This value is compatible with the peak flux observed in the last outburst detected in 2009 (Bozzo et al. 2010). Assuming a bolometric correction factor of ~ 2 (estimated from the ratio between the flux estimated in the 0.5–100 keV and 0.5–10 keV energy band; see Table 2), the observed peak flux corresponds to a 0.5–100 keV luminosity of $\simeq 8 \times 10^{36} d_{6.9}$ erg s $^{-1}$. The outburst flux decayed during ≈ 25 d before the source became no longer detectable. The spectral shape observed by *Swift*-XRT was roughly constant, with a softening of the spectrum at the end of the outburst.

Similarly to the 2009 outburst (Papitto et al. 2010; Falanga et al. 2011; Ibragimov et al. 2011), and to other AMPSs (Patruno & Watts 2012), the X-ray spectrum of IGR J17511–3057 observed during the 2015 outburst was dominated by a hard component described by a power-law with index $\Gamma \simeq 1.6\text{--}1.8$, and a cut-off at a temperature of $\gtrsim 20$ keV. The emission observed by *XMM-Newton* and *INTEGRAL* during the 2015 outburst was modelled as thermal Comptonization of soft photons with temperature $kT_s \gtrsim 1$ keV, off a cloud of hot electrons ($kT_e \gtrsim 20$ keV). Two thermal components were detected at soft X-rays, and interpreted as radiation coming from an accretion disk truncated a few tens of km from the NS surface and from the NS surface itself. Signatures of disk reflection such as a broadened ($\sigma \approx 1$ keV) line centered at an energy compatible with ionized iron was also observed in the *XMM-Newton* spectrum.

X-ray pulsations at the spin period of IGR J17511–3057 were detected by the *XMM-Newton* EPIC-pn at an RMS amplitude of $\simeq 15$ per cent in the 0.5–10 keV band, and

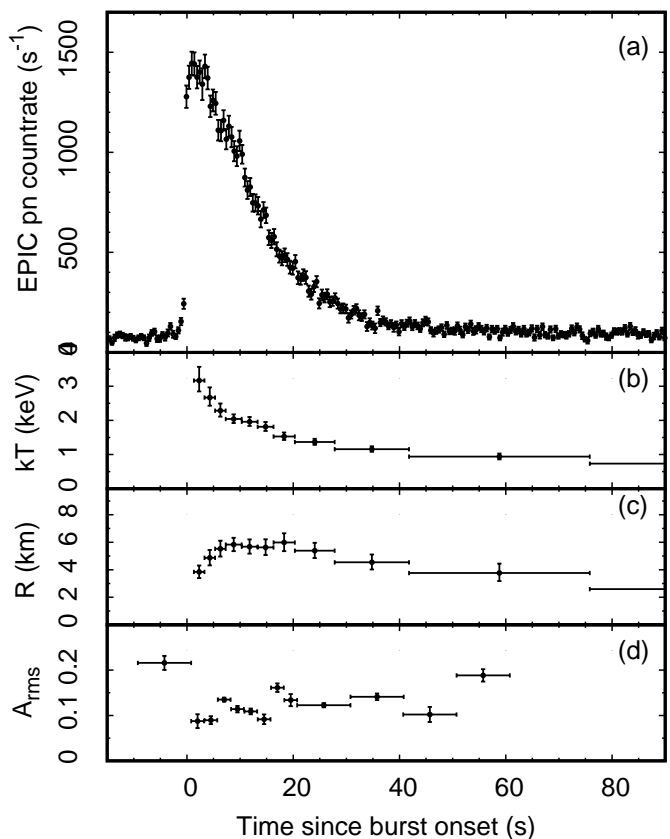


Fig. 7. 0.3–10 keV light curve of the first type-I X-ray burst observed by *XMM-Newton* during the 2015 observation (panel a). The burst onset took place at MJD 57108.01023. The temperature and the apparent radius (evaluated considering a distance to the source of 6.9 kpc) are plotted in panels (b) and (c). Panel (d) shows the rms amplitude of pulsations at the spin frequency of the pulsar (see Table 3).

by *INTEGRAL* at an amplitude of 14 per cent in the 20–100 keV band. Up to four harmonics were needed to model the pulse profile observed in the soft X-ray band, similar to the case of the 2009 outburst (Papitto et al. 2010). We searched for kHz quasi-periodic oscillations as those reported by Kalamkar et al. (2011), but did not find any hint in the Fourier power spectrum.

The relatively short exposure (~ 70 ks) over which pulsations were detected prevented us from studying the frequency evolution driven by the accretion of matter. Only a loose upper limit on the spin frequency derivative of $\sim 10^{-11}$ Hz s $^{-1}$ could be obtained. On the other hand, the expected spin up rate driven by accretion torques is:

$$\dot{\nu} \simeq \frac{1}{2\pi I_*} \frac{R_* L_X}{\sqrt{GM_*}} \sqrt{R_{in}}, \quad (2)$$

where M_* , R_* and I_* are the NS mass, radius, and moment of inertia, respectively, R_{in} is the disk truncation radius and L_X is the accretion luminosity. The accretion luminosity derived from the flux observed in the 0.5–100 keV band is $L_X = 8.1 \times 10^{36} d_{6.9}$ erg s $^{-1}$. Evaluating Eq. 2 for a disk truncated at the corotation radius ($R_c = 42.8 m_{1.4}^{1/3}$ km for IGR J17511–3057), and taking $M_* = 1.4 M_\odot$, $R_* = 10$ km and $I = 10^{45}$ g cm 2 , yields $\dot{\nu} \simeq 2 \times 10^{-13}$ Hz s $^{-1}$, two orders of magnitude lower than the upper limit we could

set based on the *XMM-Newton* observation alone. A spin-up at an average rate of $\dot{\nu}_{09} = (1.45 \pm 0.16) \times 10^{-13} \text{ Hz s}^{-1}$ was indeed reported by Riggio et al. (2011b) for the 2009 outburst, when the source emitted a luminosity similar to that observed during the 2015 outburst.

Assuming that the spin frequency observed by *XMM-Newton* during a ≈ 70 ks exposure is a good tracer of the actual spin frequency of the pulsar, we can derive constraints on the evolution of the spin frequency of the pulsar during the $\Delta t_q = 1991$ d elapsed between the end of 2009 outburst (MJD 55113; see Fig. 1 of Riggio et al. 2011b), and the beginning of the 2015 outburst (MJD 57104; Bozzo et al. 2015a). In 2009, *XMM-Newton* measured an average spin frequency of $\nu_{09} = 244.83395121(4) \text{ Hz}$ (Papitto et al. 2010), and the frequency decrease between this value and that measured in 2015 outburst is $\Delta\nu = (\nu_{15} - \nu_{09}) = (-1.0 \pm 0.5) \times 10^{-7} \text{ Hz}$. To evaluate the effect on this value of the positional uncertainty we maximized the difference between the frequency shift evaluated using Eq. 1 for the 2009 and the 2015 outburst, respectively, obtaining $\sigma_{\Delta\nu}^{pos} = 0.5 \times 10^{-7} \text{ Hz}$. Adding in quadrature this term to the statistical uncertainty yields an estimate of the frequency change $\Delta\nu = (-1.0 \pm 0.7) \times 10^{-7} \text{ Hz}$, compatible with zero within less than 2σ . A spin down rate of $\text{few} \times 10^{-15} \text{ Hz s}^{-1}$ has been observed from all the four AMSPs for which pulsations were detected during consecutive outbursts, so far (Hartman et al. 2008, 2009; Patruno et al. 2012; Patruno 2010; Hartman et al. 2011; Papitto et al. 2011; Riggio et al. 2011a; Patruno et al. 2009). It is therefore useful to quote the measured 3σ upper limit on the spin frequency decrease of IGR J17511–3057 ($|\Delta\nu| < 3 \times 10^{-7} \text{ Hz}$), and the corresponding limit on the average rate of spin frequency variation ($|\dot{\nu}| < 1.7 \times 10^{-15} \text{ Hz s}^{-1}$).

A more conservative upper limit of the magnitude of the spin frequency derivative of IGR J17511–3057 during quiescence was obtained taking into account that the pulsar spun up during the ≈ 18 d elapsed since the 2009 *XMM-Newton* observation and the end of that outburst (Riggio et al. 2011b), and assuming that it did the same during the nearly three days occurred since the 2015 outburst beginning and the *XMM-Newton* observation. Considering the spin-up rate derived by Riggio et al. (2011b) ($\dot{\nu}_{09} = (1.45 \pm 0.16) \times 10^{-13} \text{ Hz s}^{-1}$), we estimated a frequency increase of $\Delta\nu^{accr} \simeq (+2.6 \pm 0.3) \times 10^{-7} \text{ Hz}$ taking place at the end of the 2009 outburst and at the beginning of the 2015 outburst. Adding this value to the upper limit on the measured frequency change, we obtained an upper limit on the spin-down rate of the pulsar during quiescence of $|\Delta\nu| < |\Delta\nu^{accr} - \Delta\nu| < 6 \times 10^{-7} \text{ Hz}$. The corresponding limit on the spin down rate is $|\dot{\nu}| < 3.5 \times 10^{-15} \text{ Hz s}^{-1}$.

Assuming that the main torque operating on AMSPs during quiescence was due to electromagnetic spin-down, the upper limit on the spin down rate can be translated into a constraint on the strength of the magnetic dipole moment. Considering the relation derived by Spitkovsky (2006) for the spin-down of a force-free, rotating magnetosphere, one can estimate the NS magnetic dipole moment, μ , needed to produce a certain spin down as

$$\mu = \left[\frac{I_* c^3 \dot{\nu}}{(1 + \sin^2 \alpha)(2\pi)^2 \nu^3} \right]^{1/2}. \quad (3)$$

Here, α is the magnetic inclination angle. Setting $\alpha = 30^\circ$, and considering the conservative upper limit on the spin

down rate estimated earlier ($|\dot{\nu}| < 3.5 \times 10^{-15} \text{ Hz s}^{-1}$) yielded an upper limit of $\mu < 3.6 \times 10^{26} \text{ G cm}^3$.

In principle, the torque acting on a pulsar that is propelling away the inflowing matter through centrifugal inhibition of accretion (the so-called propeller effect, Illarionov & Sunyaev 1975) is an alternative explanation of the spin down of AMSP during quiescence. Assuming that an accretion disk surrounds the pulsar even during quiescence and is truncated at a radius R_{in} exceeding the co-rotation radius R_c , the propeller spin down torque is $N_{sd} = \dot{M}_d \sqrt{GM R_{in}} [(R_{in}/R_c)^{2/3} - 1]$ (see, e.g., Eq. 17 of Papitto & Torres 2015, evaluated for completely anelastic propeller interaction). Assuming $R_{in} = 80 \text{ km}$, a disk mass accretion rate of $\approx 1.5 \times 10^{-11} \text{ M}_\odot \text{ yr}^{-1}$ is required to explain a spin-down rate of $\simeq 3 \times 10^{-15} \text{ Hz s}^{-1}$ in terms of the propeller effect. The propeller luminosity associated to such a value is $L_{prop} \simeq 2.5 \times 10^{34} \text{ erg s}^{-1}$ (see Eq. 18 of Papitto & Torres 2015). No estimates of the high energy emission during quiescence are available for IGR J17511–3057. However, such a value exceeds by up to two orders of magnitude the X-ray emission observed from other AMSPs in quiescence, and we deem it as unlikely. A value of the X-ray luminosity in excess of $\approx 10^{33} \text{ erg s}^{-1}$ would then be taken as an indication of the presence of strong outflows from the system, as it was done by Papitto et al. (2014); Papitto & Torres (2015) to interpret the emission of two transitional millisecond pulsars in the accretion disk state.

We detected three consecutive type-I X-ray bursts during the *XMM-Newton* observation performed in 2015. Two more were detected by *INTEGRAL* but the much larger time elapsed between them ($\approx 3.5 \text{ d}$) indicates that they were not consecutive. No evidence of burst radius expansion was detected in the three bursts. The peak flux observed during the bursts presented here ($\sim 1.5 \times 10^{-8} \text{ erg cm}^{-2} \text{ s}^{-1}$ in the 0.5–10 keV band, $\sim 3 \times 10^{-8} \text{ erg cm}^{-2} \text{ s}^{-1}$ in the 3–20 keV) is not larger than the value reported by Altamirano et al. (2010) for the burst seen during the 2009 outburst ($6.7 \times 10^{-8} \text{ erg cm}^{-2} \text{ s}^{-1}$). As a consequence, the non observation of photospheric radius expansion during the bursts analyzed here could not yield a tighter constraint on the source distance than that reported by Altamirano et al. (2010, 6.9 kpc). During the 2009 event, eighteen bursts were detected, each characterized by an exponential decay timescale of $\simeq 10 \text{ s}$ (Altamirano et al. 2010; Bozzo et al. 2010; Papitto et al. 2010; Falanga et al. 2011). The dependence of the burst recurrence time t_{rec} on the persistent X-ray flux was studied by Falanga et al. (2011), who found a dependence $t_{rec} \propto F_{pers}^{-1.1}$. Using the relation plotted in their Fig. 10, at a bolometric flux of $\simeq 1.4 \times 10^{-9} \text{ erg cm}^{-2} \text{ s}^{-1}$ (such as that deduced from the simultaneous *XMM-Newton-INTTEGRAL* spectral modelling, see Table 2), a recurrence time of $\simeq 9 \text{ hr}$ is expected. This is only slightly larger than the recurrence time observed by *XMM-Newton* in 2015 ($\simeq 8 \text{ hr}$).

5. Conclusions

We have presented an analysis of *XMM-Newton*, *INTEGRAL* and *Swift* observations performed during the outburst detected from IGR J17511–3057 during early 2015, the second observed from the pulsar after the discovery outburst in 2009. The outburst profile, spectral and burst properties were remarkably similar to those observed during the last accretion event detected in 2009, suggesting that the

properties of the accretion flow did not change much in the two episodes. The frequency of the coherent signal detected by *XMM-Newton* three days after the beginning of the outburst was compatible with the value measured by the same observatory during the 2009 outburst. Therefore, a firm assessment of the spin evolution of the pulsar during the time elapsed between the two outbursts was not possible. However, taking into account the accretion driven spin up observed during the 2009 outburst, and assumed for the 2015 outburst, we derived an upper limit of $3.5 \times 10^{-15} \text{ Hz s}^{-1}$ on the spin down rate during quiescence. Electromagnetic spin-down of a NS with a magnetic field weaker than $3.5 \times 10^8 \text{ G}$ (at the equator of the NS and assuming a magnetic inclination of 30°) can account for this inferred spin-down. A magnetic field of the same order has been inferred also for other AMSPs. Observations of future outbursts will allow to derive more stringent constraints on the long-term spin evolution of the pulsar, as well as enabling a first estimate of the orbital period derivative.

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